

$$E = mc^2$$
$$= \hbar\omega$$

Holger Müller

Kioloa 2011

Compton oscillators + time dilation = de Broglie (1924) relations



$$h\nu_0 = m_0c^2$$

ν_0 being measured, of course, in a system which is at rest with respect to a certain amount of energy. This hypothesis is the base of our system, and it is

Then the frequency observed by the fixed observer will be

$$\nu_1 = \nu_0(1 - \beta^2)^{1/2} = \frac{m_0c^2}{h}(1 - \beta^2)^{1/2}$$

On the other hand, since the energy of the moving body is equal to $m_0c^2/(1 - \beta^2)^{1/2}$ for the same observer, the corresponding frequency according to the relationship of the quantum is

$$\nu = h^{-1}[m_0c^2/(1 - \beta^2)^{1/2}].$$

The two frequencies ν_1 and ν are essentially different since the factor $(1 - \beta^2)^{1/2}$ is not involved in the same way. This is a difficulty that has intrigued me for a long time; I have succeeded in eliminating it by demonstrating the following theorem that I shall call the theorem of the

De Broglie + Gravity

Particle represents an oscillator

$$\omega_C = \frac{mc^2}{\hbar}.$$

Time dilation

$$\omega_C \rightarrow \omega_C \sqrt{1 - v^2 / c^2}$$

Gravitational redshift

$$\omega_C \rightarrow \omega_C \left(1 - \frac{\Delta U}{c^2} \right)$$

Time dilation + redshift

$$\omega_C \rightarrow \omega_C \sqrt{-g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}}$$

Proper time $\sim$ action

Proper time $c d\tau = \sqrt{-g_{\alpha\beta} dx^\alpha dx^\beta}$

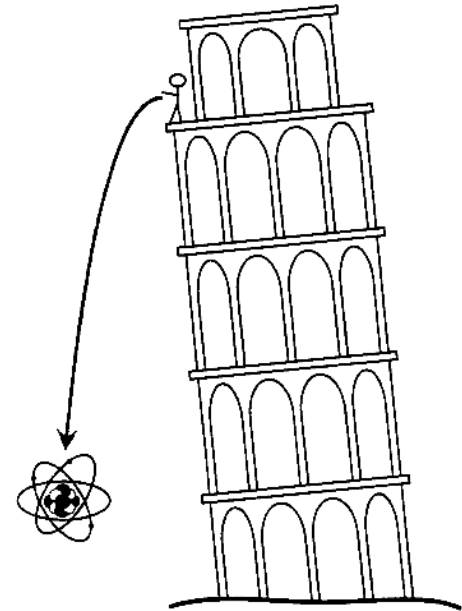
Action $S = mc^2 \int d\tau$

Path integral $\Psi(t^B; \vec{x}^B) = \int Dq \exp\left(\frac{i}{\hbar} S\right) \Psi(t^A; \vec{x}^A)$

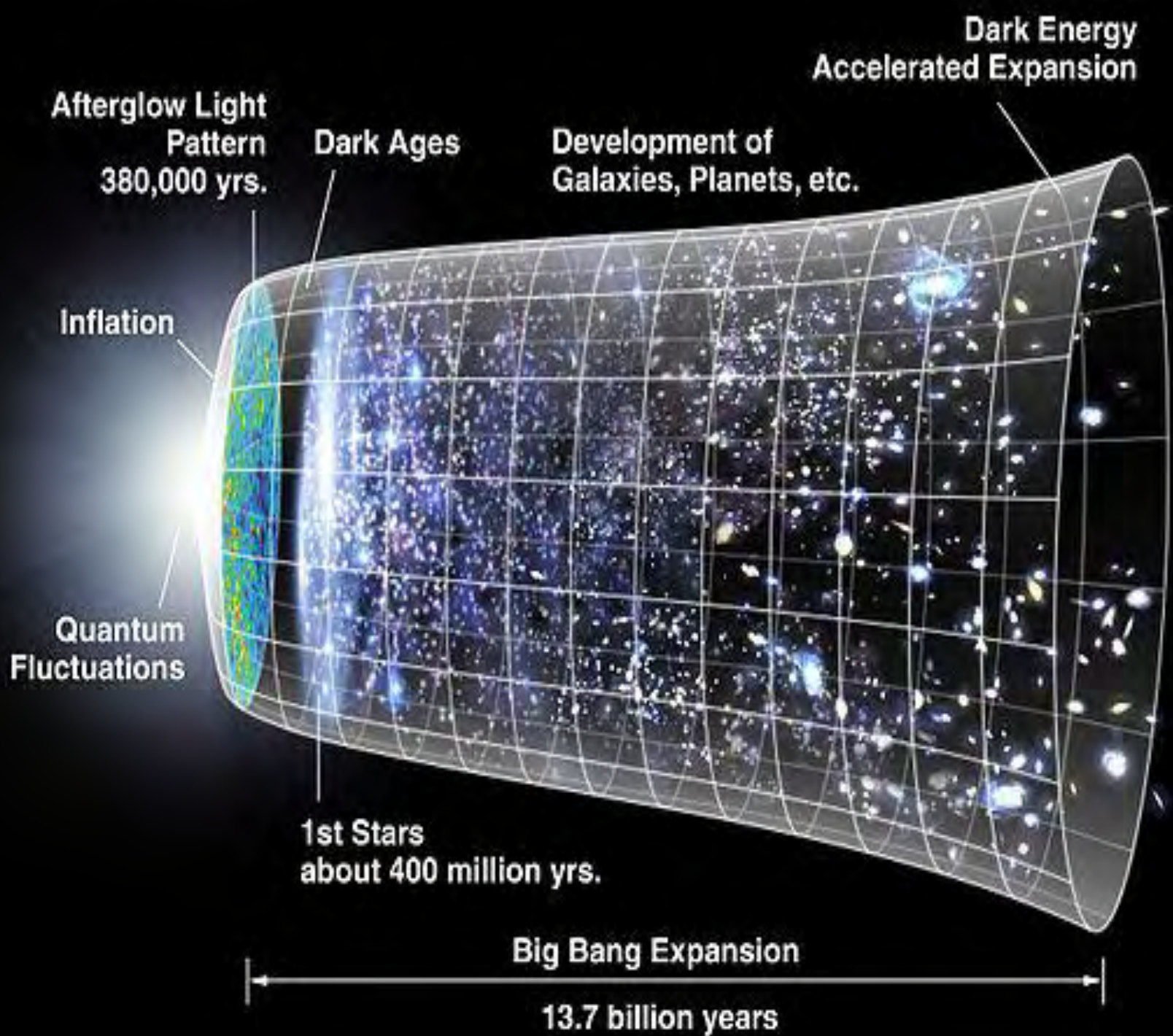
Compton oscillators + general relativity = quantum mechanics

Outline

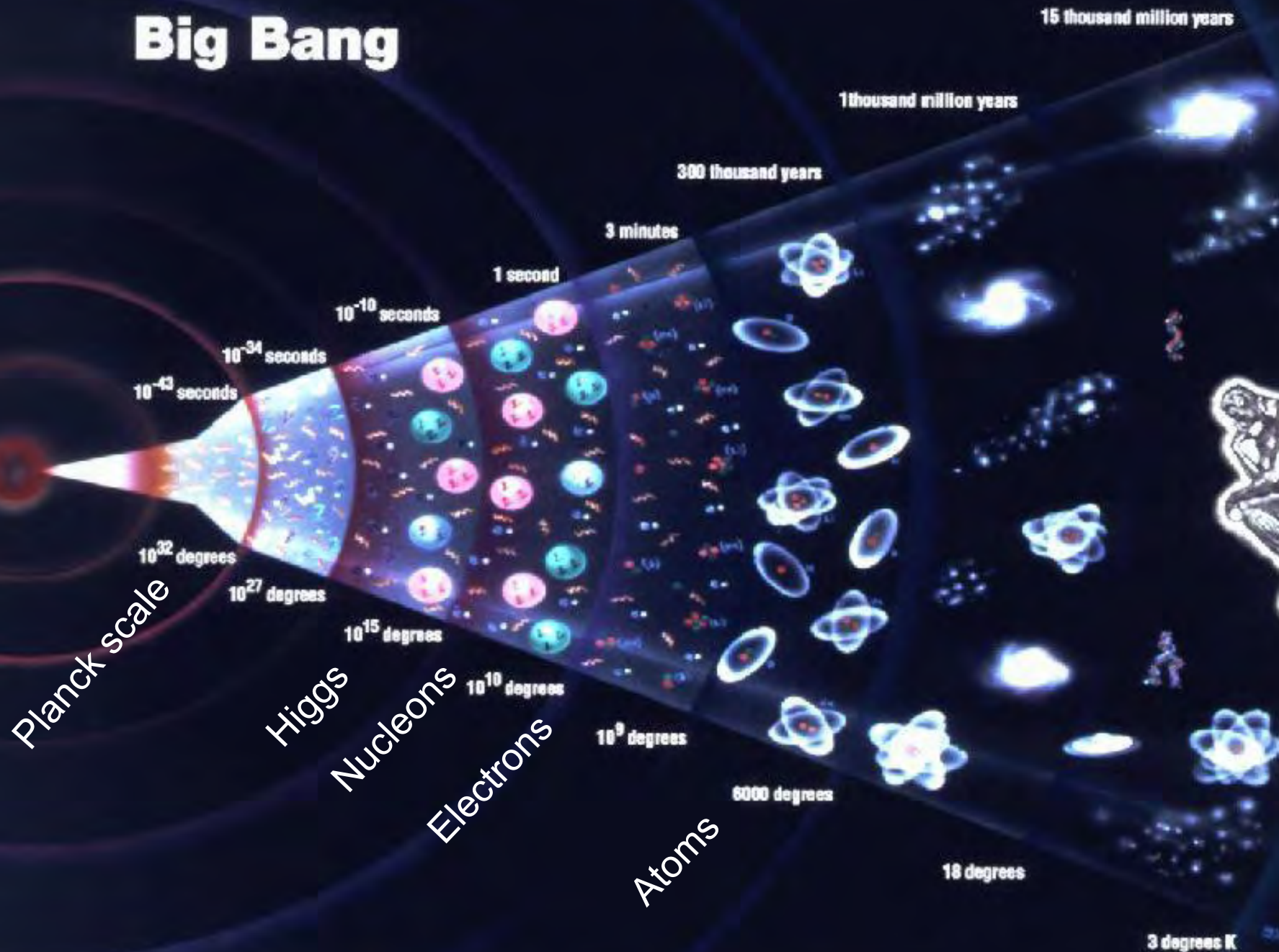
1. Motivation
2. Atom interferometer redshift test
3. Really? Don't you just measure g , not U ?
4. Aharonov-Bohm effect
 - we measure U
5. Compton clock
 - particles are *really* clocks
6. What is time?



A. Peters



Big Bang



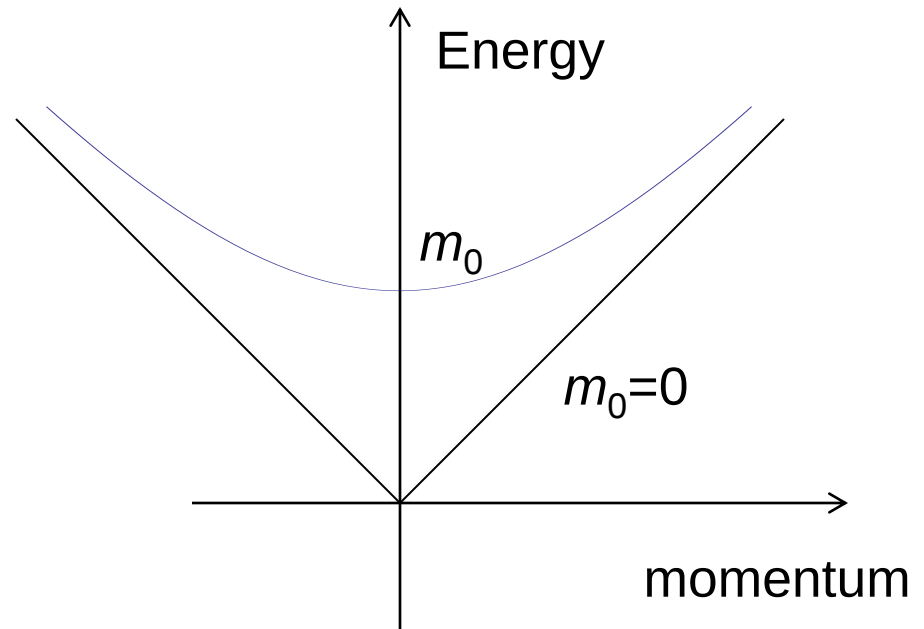
Penrose's argument

If you have nothing that has rest mass, you can't have a clock. And you can't measure distances either.

Maxwell equations are scale invariant

- Beginning: All $m=0$.
- Big freeze: only photons left. Again, $m=0$

“The scale changed,” but there is nothing to measure scales



Is a rock a clock?

“Ticking” of an atomic ($^{27}\text{Al}^+$) clock

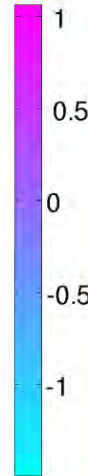
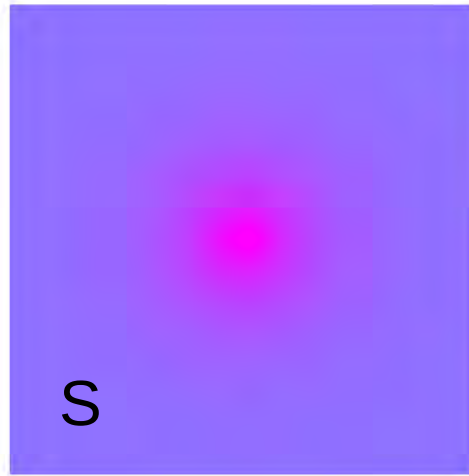
0.000 000 0045 GeV



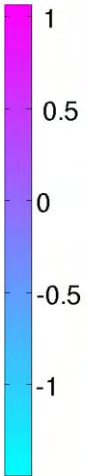
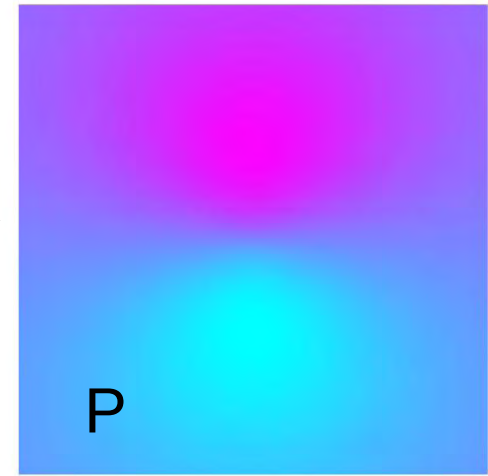
27 GeV



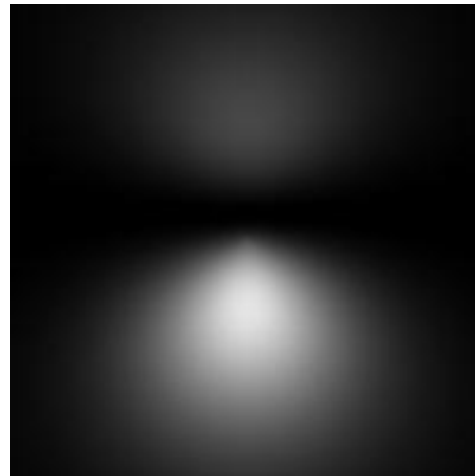
$E=0$



$+e^{-i\omega t}$



=



$\omega = 2\pi \times 1.121 \times 10^{15}$
Hz

Thanks to Till Rosenband!

$E = mc^2 = h\nu$ provides a clock

Action

$$S = \int mc^2 d\tau,$$

Propagator

$$\langle \psi(t_1) | \psi(t_2) \rangle = \int Dq \exp(iS / \hbar).$$

Quantum phase

$$\varphi = \frac{i}{\hbar} \int mc^2 d\tau = \int \omega_C d\tau,$$

$$d\tau / dt = \sqrt{-g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu}.$$

Quantum phase proportional to proper time (!)

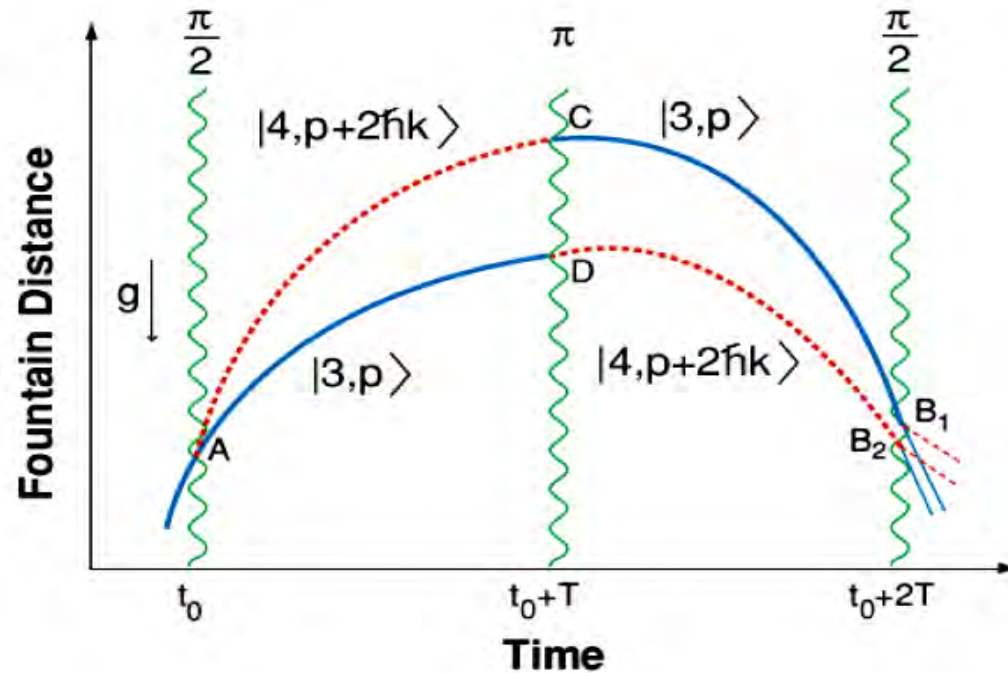
[Misner, Thorne & Wheeler (1970); Feynman & Hibbs, (1965);
Borde, *Eur. Phys. J. Special Topics* (2008).]





Atom interferometer redshift test

Atom Interferometer



$$\Delta\varphi = \frac{mc^2}{\hbar} \int \left(\underbrace{\frac{\varphi_s - \varphi_e}{c^2}}_{\substack{\text{Redshift} \\ kgT^2}} - \underbrace{\frac{\vec{v}_s^2 - \vec{v}_e^2}{2c^2}}_{\substack{\text{Time dilation} \\ -kgT^2}} \right) dt + \sum \underbrace{\varphi_{laser}}_{\substack{\text{Laser interaction} \\ kgT^2}}$$

related by $v=dx/dt$

History

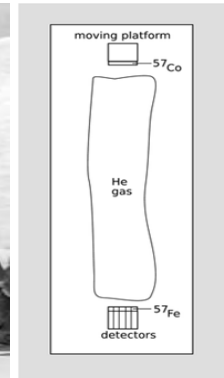
Einstein¹
1911:
solar
spectral
lines.
Didn't
work.

Einfluß der Schwerkraft auf die Ausbreitung des Lichtes. 905

müssen also die Spektrallinien des Sonnenlichtes gegenüber den entsprechenden Spektrallinien irdischer Lichtquellen etwas nach dem Rot verschoben sein, und zwar um den relativen Betrag

$$\frac{\nu_0 - \nu}{\nu_0} = \frac{-\Phi}{c^2} = 2 \cdot 10^{-6}.$$

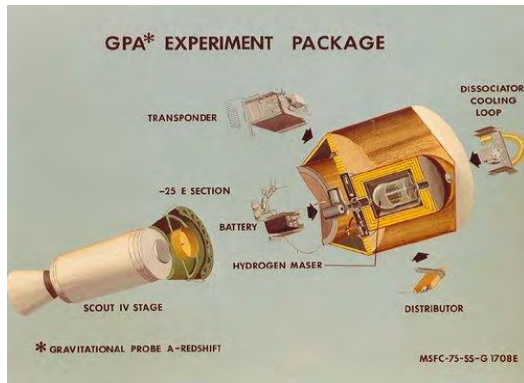
Wenn die Bedingungen, unter welchen die Sonnenlinien entstehen, genau bekannt wären, wäre diese Verschiebung noch der Messung zugänglich. Da aber anderweitige Einflüsse (Druck, Temperatur) die Lage des Schwerpunktes der Spektrallinien beeinflussen, ist es schwer zu konstatieren, ob der hier



Pound & Rebka,²
1960:
Mössbauer effect.

10⁻².

Vessot *et al.*,⁴
1976:
Atomic
clocks in
rocket
10⁻⁵.



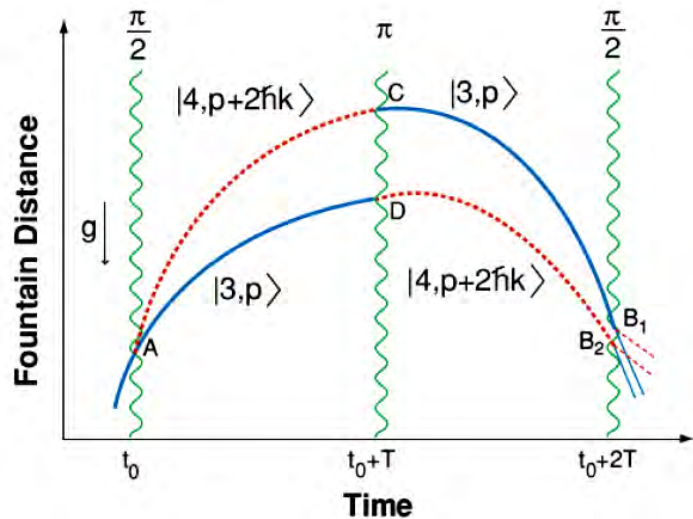
Hafele & Keating,³
1971:
Atomic
clocks in
aircraft

10⁻¹.

¹ Ann. Phys. **340**, 898 (1911). ² PRL **4**, 337 (1960); **13**, 539 (1964); Phys. Rev. **140**, B788 (1965).

³ Science **177**, 166 (1972); 177, 168 (1972). ⁴ PRL **45**, 2081 (1980).

Atom Interferometer



Redshift Time dilation Laser

$$\Delta\varphi = \frac{mc^2}{\hbar} \int \left(\underbrace{\frac{\varphi_s - \varphi_e}{c^2} - \frac{\vec{v}_s^2 - \vec{v}_e^2}{2c^2}}_{=0} \right) dt + \sum_{i=1}^3 k_i z(t_i)$$

(non-gravitational energy conservation)

Alternative Interpretation

- If *any* energy is conserved, first terms also cancel
- then, testing UFF = testing redshift
- AI becomes indirect measurement of g
- As any redshift measurement

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A precision measurement of the gravitational redshift by the interference of matter waves

Holger Müller^{1,2}, Achim Peters³ & Steven Chu^{1,2,4}

One of the central predictions of metric theories of gravity, such as general relativity, is that a clock in a gravitational potential U will run more slowly by a factor of $1 + U/c^2$, where c is the velocity of light, as compared to a similar clock outside the potential¹. This effect, known as gravitational redshift, is important to the operation of the global positioning system², timekeeping^{3,4} and future experiments with ultra-precise, space-based clocks⁵ (such as searches for variations in fundamental constants). The gravitational redshift has been measured using clocks on a tower⁶, an aircraft⁷ and a rocket⁸, currently reaching an accuracy of 7×10^{-5} . Here we show that laboratory experiments based on quantum interference of atoms^{9,10} enable a much more precise measurement, yielding an accuracy of 7×10^{-9} . Our result supports the view that gravity is a manifestation of space-time curvature, an underlying principle of general relativity that has come under scrutiny in connection with the search for a theory of quantum gravity¹¹. Improving the redshift measurement is particularly important because this test has been the least accurate among the experiments that are required to support curved space-time theories¹.

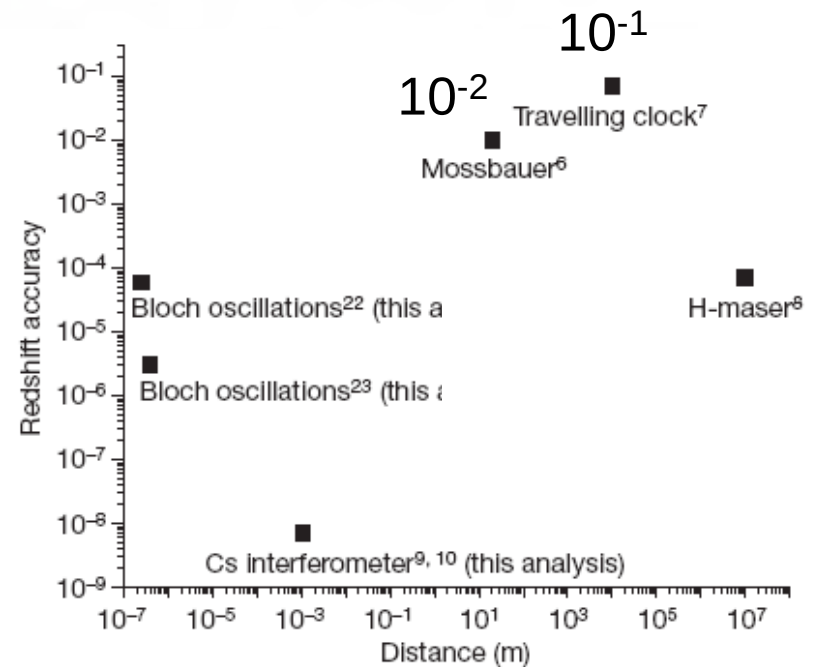


Figure 2 | Absolute determinations of the gravitational redshift. The accuracy (defined as the standard error) in β is plotted versus the relative height of the clocks.

Criticism

- ω_C is unphysical: the AI is not a pair of clocks, but just a measurement of g
- The atom is no real clock: it cannot be used to lock an oscillator.

Does an atom interferometer test the gravitational redshift at the Compton frequency ?

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Reynaud,⁴ Christophe Salomon,⁵ and Claude Cohen-Tannoudji⁵

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ENS, UPMC, 24 rue Lhomond, 75231 Paris, France*

(Dated: May 13, 2011)

Abstract

Atom interferometers allow the measurement of the acceleration of freely falling atoms with respect to an experimental platform at rest on Earth's surface. Such experiments have been used



Equivalence Principle and Gravitational Redshift

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(Received 17 February 2011; published 11 April 2011)

We investigate leading order deviations from general relativity that violate the Einstein equivalence principle in the gravitational standard model extension. We show that redshift experiments based on matter waves and clock comparisons are equivalent to one another. Consideration of torsion balance tests, along with matter-wave, microwave, optical, and Mössbauer clock tests, yields comprehensive limits on spin-independent Einstein equivalence principle-violating standard model extension terms at the 10^{-6} level.

DOI: [10.1103/PhysRevLett.106.151102](https://doi.org/10.1103/PhysRevLett.106.151102)

PACS numbers: 04.80.-y, 03.30.+p, 11.30.Cp, 12.60.-i

TABLE I. Sensitivity of redshift experiments. The EEP-violation signal for each experiment is given as a linear combination of SME parameters. The observable for the Pound-Rebka Mössbauer test, e.g., is $-1.1 \text{ GeV}^{-1} \alpha(\bar{a}_{\text{eff}}^n)_0 - 1.1 \text{ GeV}^{-1} \alpha(\bar{a}_{\text{eff}}^{e+p})_0 + (-0.34 + [-0.66])(\bar{c}^n)_{00} + (-0.34 + [-0.006])(\bar{c}^p)_{00} + 0.0002(\bar{c}^e)_{00}$, with $\bar{a}_{\text{eff}}^{e+p} = \bar{a}_{\text{eff}}^p + \bar{a}_{\text{eff}}^e$. The last column shows the measured value and 1σ uncertainty. Signals dependent on models for ξ are in square brackets. Curly brackets mark expected limits.

Method	$\alpha(\bar{a}_{\text{eff}}^n)_0$ GeV	$\alpha(\bar{a}_{\text{eff}}^{e+p})_0$ GeV	$(\bar{c}^n)_{00}$	$(\bar{c}^p)_{00}$	$(\bar{c}^e)_{00}$	Limit ppm
Mössbauer effect [2]	-1.072	-1.072	$0.3358 - [2/3]$	$-0.3353 - [0.006]$	0.000 182 6	1000 ± 7600
H maser on rocket [3]	-1.072	-1.072	0.3358	$0.3353 - [0.67]$	$0.000 182 6 - [1.3]$	2.5 ± 70
Cs fountain (<i>proj.</i>) [16]	-1.072	-1.072	$0.3358 + [0.40]$	$0.34 + [0.28]$	$0.000 182 6 - [1.3]$	{2}
Bloch oscillations [4,17]	0.1632	-0.1580	$-0.051 12 - [0.0005]$	$0.049 40 + [0.0010]$	0.000 026 90	3 ± 1
Bloch oscillations [6]	0.1492	-0.1439	$-0.046 73 - [0.0006]$	$0.045 00 + [0.0008]$	0.000 024 51	0.16 ± 0.14
Cs interferometer [4]	0.1881	-0.1835	$-0.058 90 - [0.0004]$	$0.057 39 + [0.001]$	0.000 031 26	0.007 ± 0.007
Rb interferometer [18]	0.1632	-0.1580	$-0.051 12 - [0.0005]$	$0.049 40 + [0.001]$	0.000 026 90	-0.004 ± 0.007

Bottom line

- Comprehensive limits on all 5 coefficients giving EEP-violations (ppm)

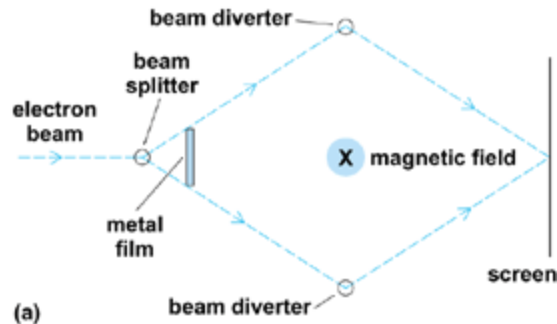
	$(a_{\text{eff}}^n)_0$	$(a_{\text{eff}}^p)_0 + (a_{\text{eff}}^n)_0$	$(c_{\text{eff}}^n)_{00}$	$(c_{\text{eff}}^p)_{00}$	$(c_{\text{eff}}^e)_{00}$
	GeV	GeV			
AI+clocks	4.3 ± 3.7	0.8 ± 1.0	7.6 ± 6.7	-3.3 ± 3.5	4.6 ± 4.6

- A sixth coefficient, $(a_{\text{eff}}^p)_0 + (a_{\text{eff}}^n)_0$, is only relevant for charged particles
- Neglected anomalies coupling to pions and mass dimension 5 operators.
- Previously, only one limit known: $|0.8(a_{\text{eff}}^n)_0 + (a_{\text{eff}}^p)_0 + (a_{\text{eff}}^n)_0| < 1 \times 10^{-11} \text{ GeV}$

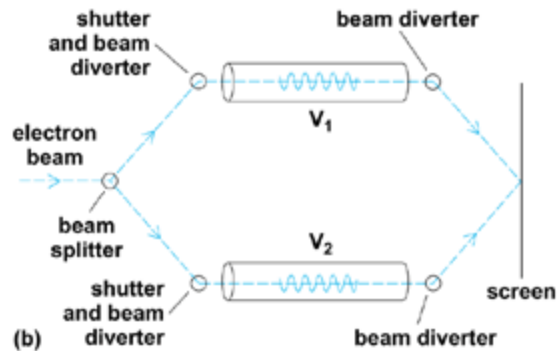


Gravitational Aharonov-Bohm Effect

Aharonov-Bohm effect



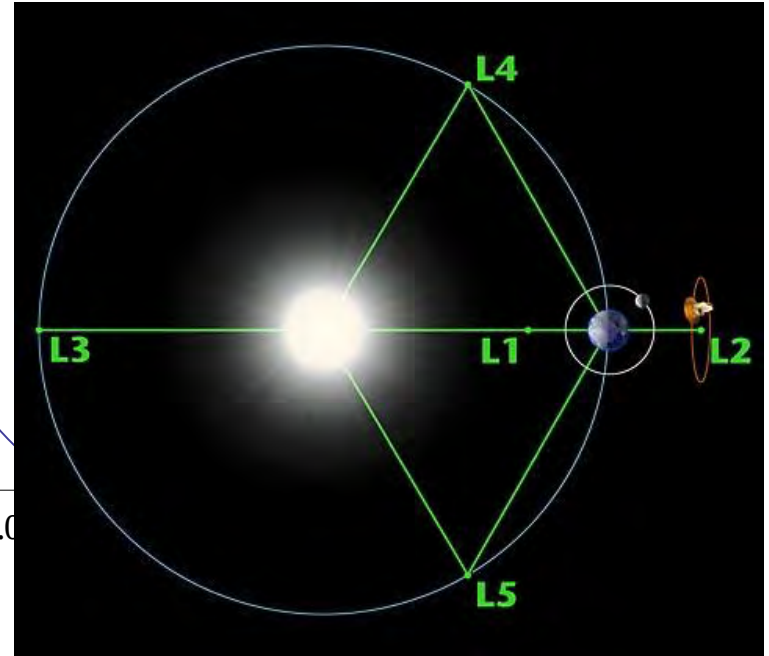
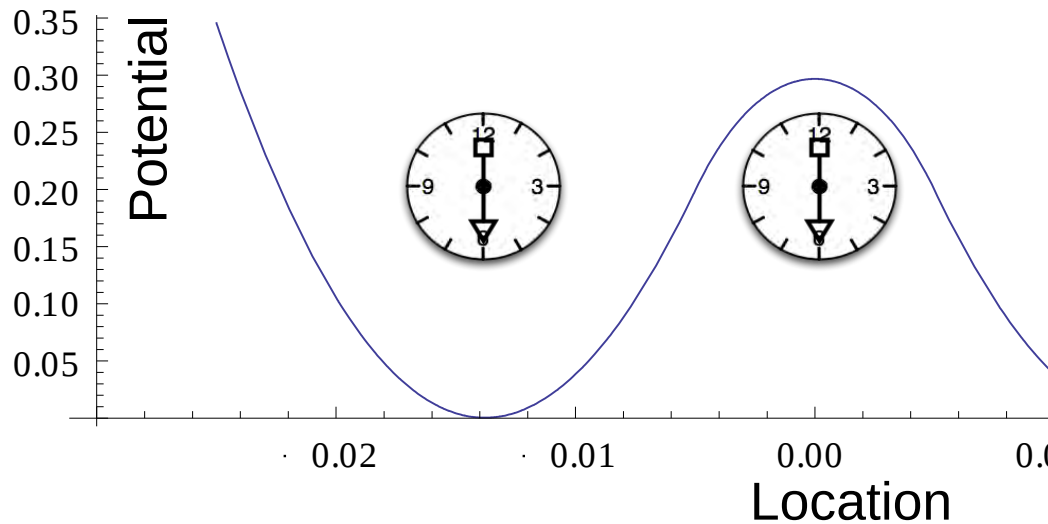
$$\varphi_A = -\frac{e}{\hbar} \int \vec{A} \cdot d\vec{l}$$



$$\varphi_V = \frac{e}{\hbar} \int V dt,$$

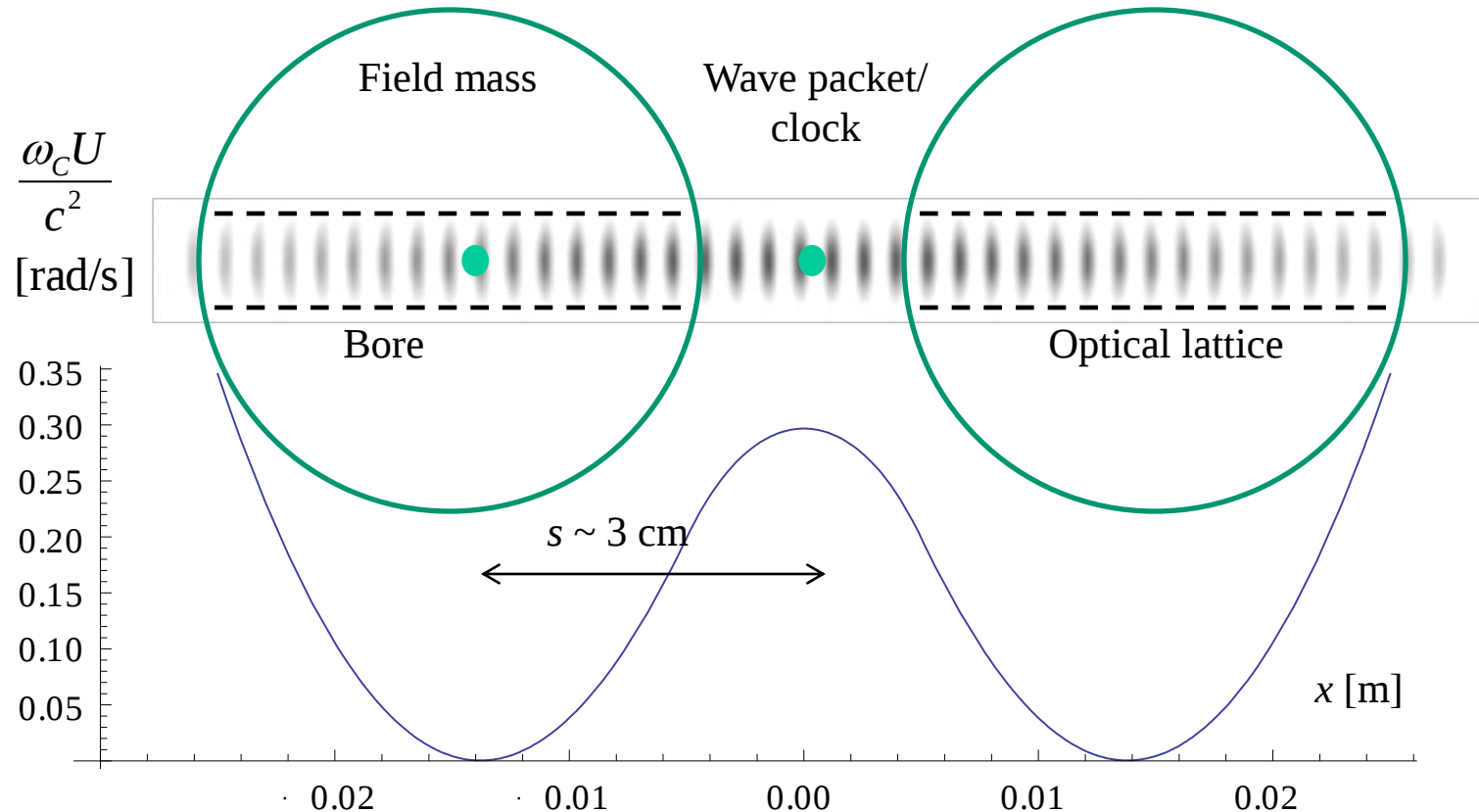
- No classical force
- Non-dispersive: not related to displacement or distortion of the wave packet
- **Topological:** cannot be predicted from any number of local measurements anywhere on the path

Gravity's Aharonov-Bohm effect



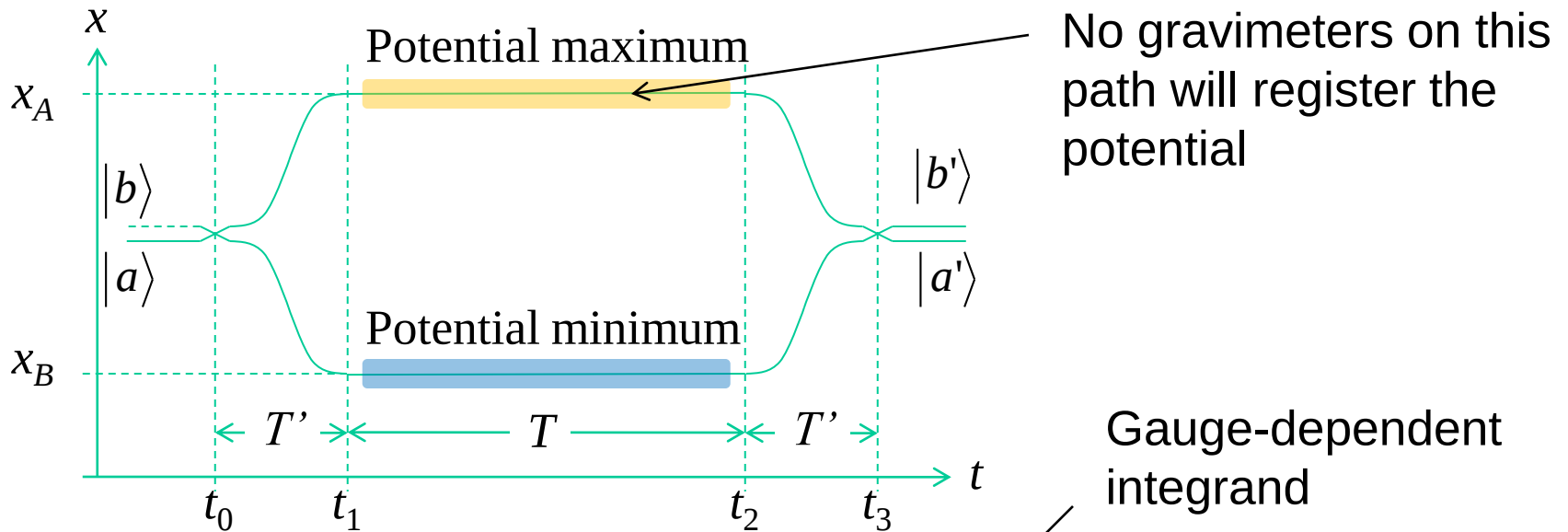
- No forces at saddle points
- Still causes redshift $\Rightarrow$ phase shift between clocks
- Terrestrial experiment: $\rho=10 \text{ g/cm}^3$, $R=10 \text{ m}$: $\Delta\nu/\nu=5 \times 10^{-21}$
- Possible realization: Earth-moon Lagrange points

Realization with atom-clocks



$$\varphi = \omega_C \int \frac{\Delta U}{c^2} dt = 0.16 \left(\frac{s}{\text{cm}} \right)^2 \left(\frac{\rho}{10 \text{ g/cm}^3} \right) \left(\frac{m}{m_{Cs}} \right) \left(\frac{T}{\text{s}} \right)$$

Properties



$$\Delta\varphi = \omega_C \oint \frac{U}{c^2} dt$$

- Topological
- Non-dispersive
- Phase shift is a consequence of the potential, not the force
- “Shaking” the path can prove time dilation as well.
- Both relativistic effects on time apply to atom-clocks
- AI is not a measurement of g

Derivation

“Quantum” picture

Time evolution

$$|\psi_{A,B}(t_2)\rangle = e^{\frac{i}{\hbar} H_{A,B} T} |\psi_{A,B}(t_1)\rangle.$$

(no field-generating masses)

$$|\psi_{A,B}(t_2)\rangle = e^{\frac{i}{\hbar} [H_{A,B} T + m U(x_{A,B}) T]} |\psi_{A,B}(t_1)\rangle.$$

(with field masses)

$$\phi_G = m \Delta U T / \hbar.$$

(differential signal)

General relativity picture

Proper time

$$\Delta \tau_0 = \frac{1}{c^2} \int [U_0(x_A) - U_0(x_B)] dt,$$

$$\Delta \phi_G = \omega \Delta \tau_G = \omega \Delta U T / c^2,$$

(differential signal)

The phase shifts are identical for clocks of $\omega_C = mc^2 / \hbar$

Systematic effects

- Residual force: displacement, not acceleration $\delta x = F/(2k^2V_0)$

Causes potential change, thus phase shift $F^2T/(4k^2V_0\hbar)$

E.g., Earth's gravity $\overline{8gGm_{Cs}^2\pi R^3\rho T}/(3\hbar k^2L^2V_0) \simeq 2 \times 10^{-9}(T/s)$

- Curvature of potential

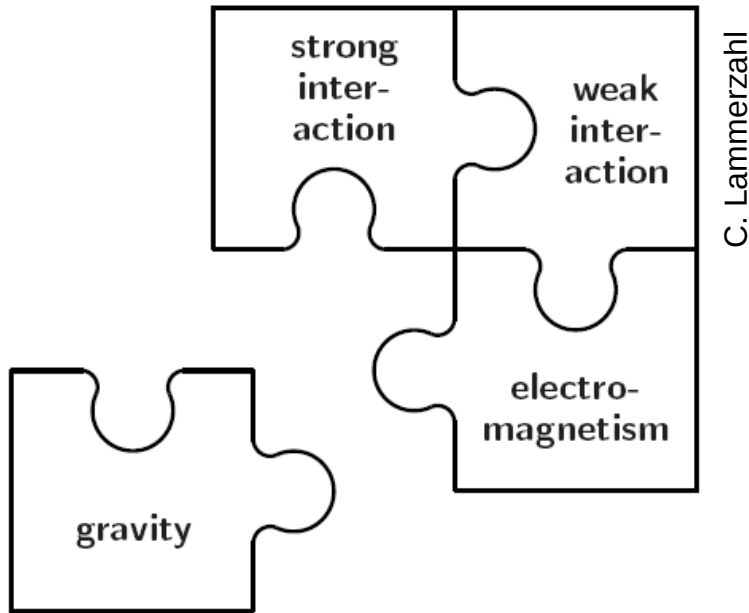
Atoms in harmonic oscillator states

Lattice
↓
 $\omega_i^2 = (\partial^2 \tilde{V} / \partial x_i^2) / (2\tilde{m})$

Curvature of gravitational potential: $V \rightarrow U + V$ $\Delta E_h = \sum_i \frac{\hbar}{2\omega_i} \partial^2 U / (\partial x_i)^2$

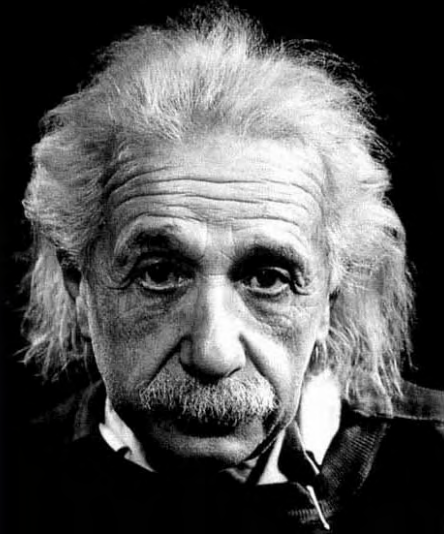
Changes

$$(2/3)\pi G\rho/\omega_i \sim 2 \times 10^{-6} \text{ rad/s},$$



“Everything should be made as simple as possible, but not simpler.”

Albert Einstein



General Relativity:

- Theory of gravity, space & time used in GPS, astrophysics, cosmology, string theory,...
- Nonlinear theory: gravity causes more gravity
- Cannot be quantized
- Low-energy limit of string theory might be Lorentz-violating [V.A. Kostelecky ...]

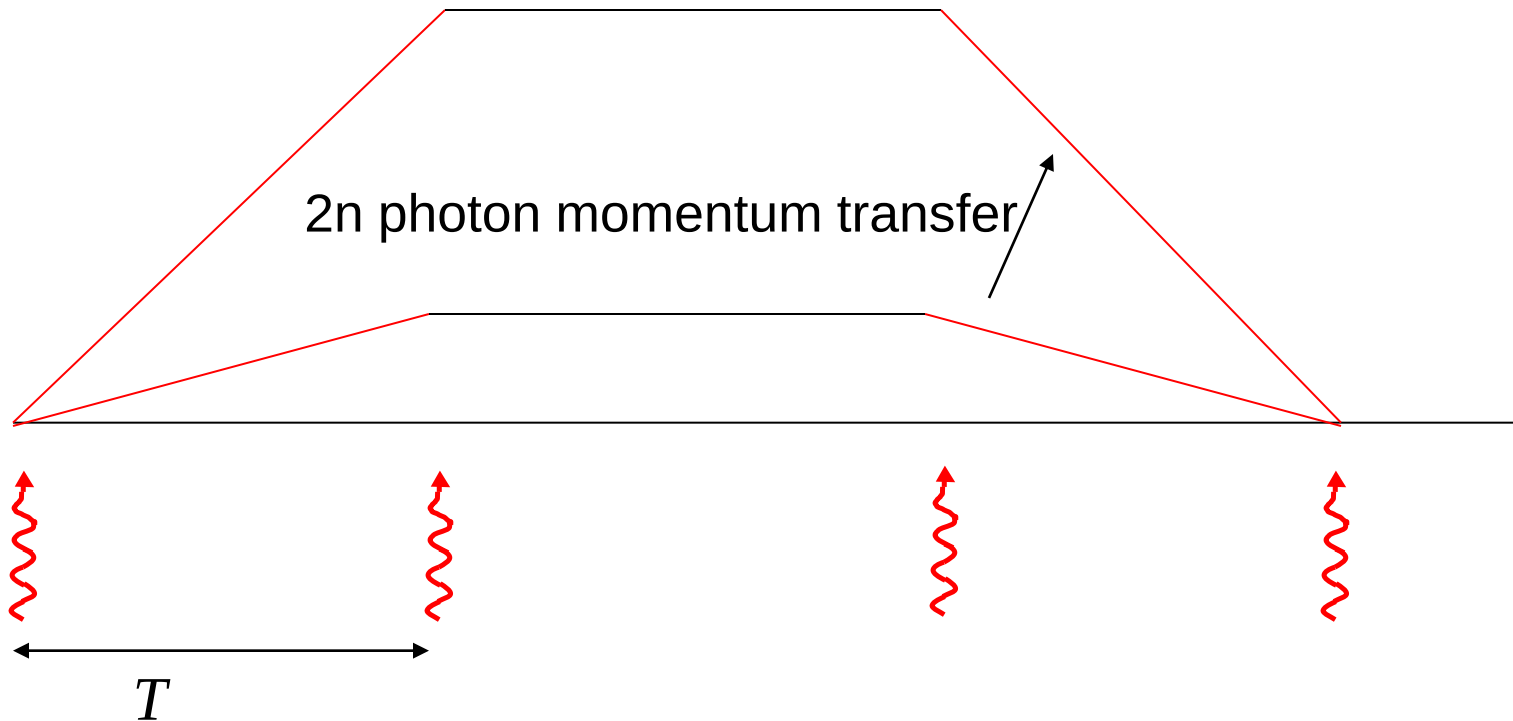


Compton clock

$$\omega = \frac{mc^2}{\hbar} \frac{1}{n^2}$$



Ramsey-Borde interferometer

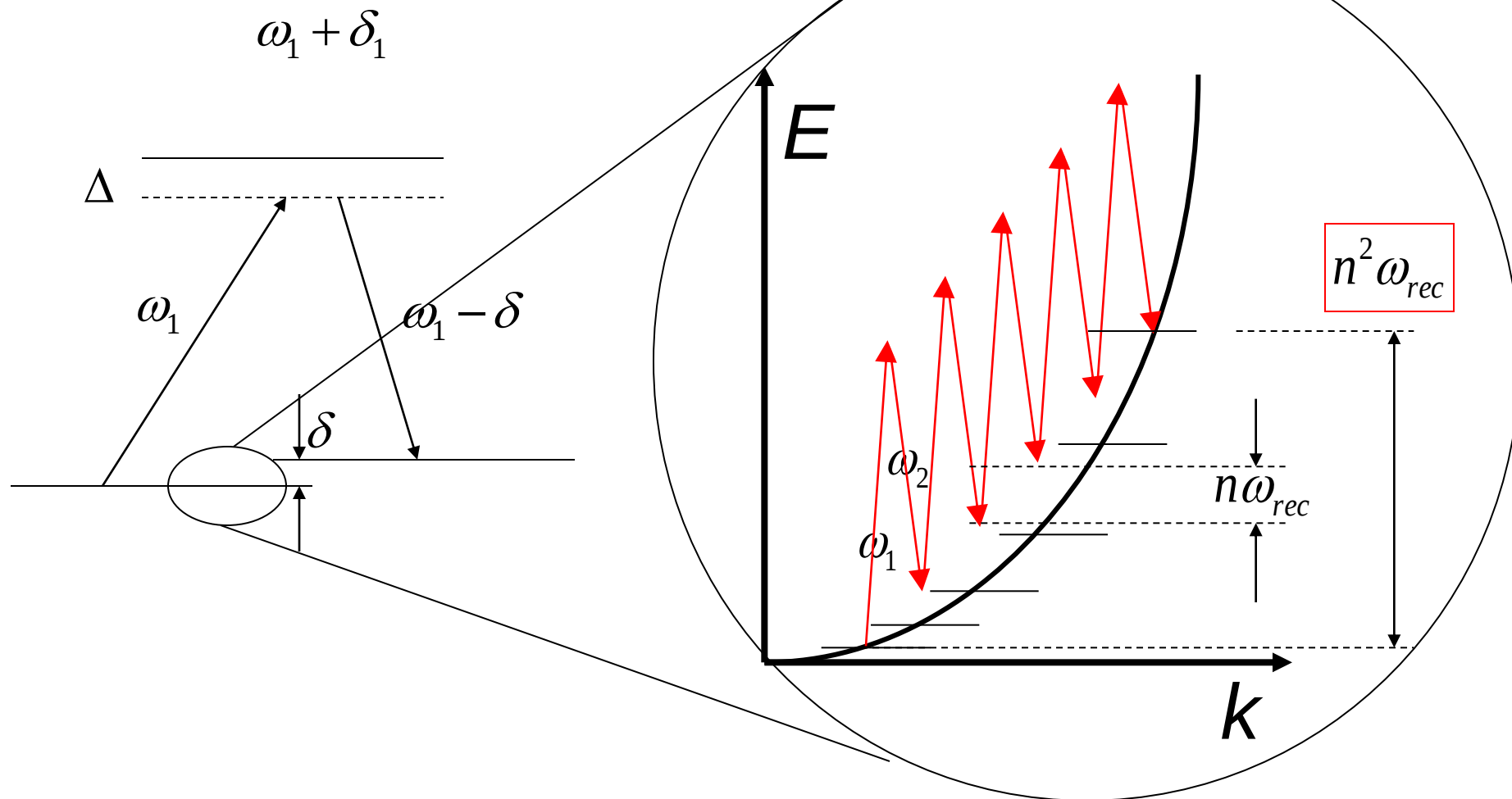


$$\Phi_1 - \Phi_2 = 16n^2 \omega_r T$$

$$\omega_r = \frac{\hbar k^2}{2m}$$

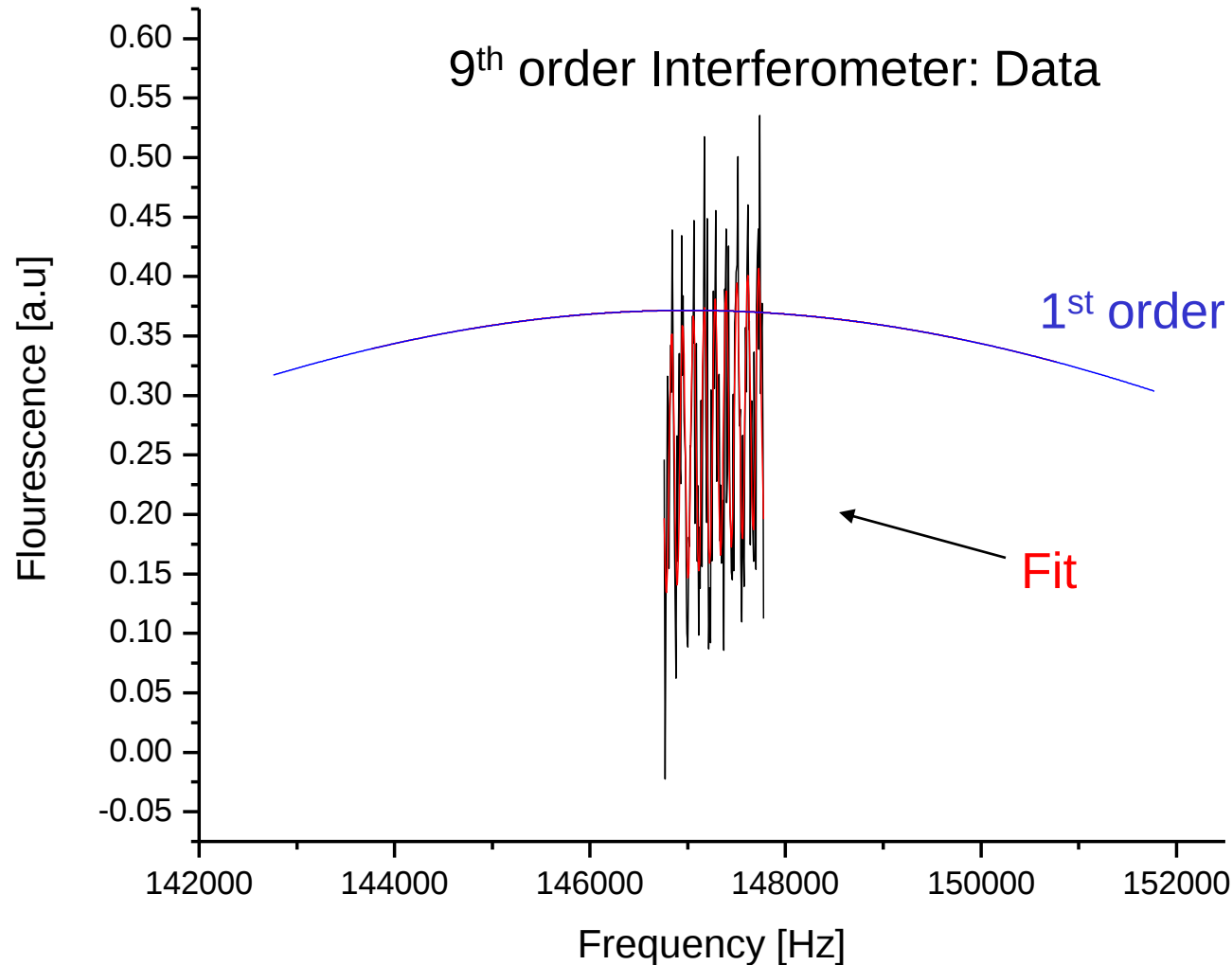


Multiphoton Bragg diffraction



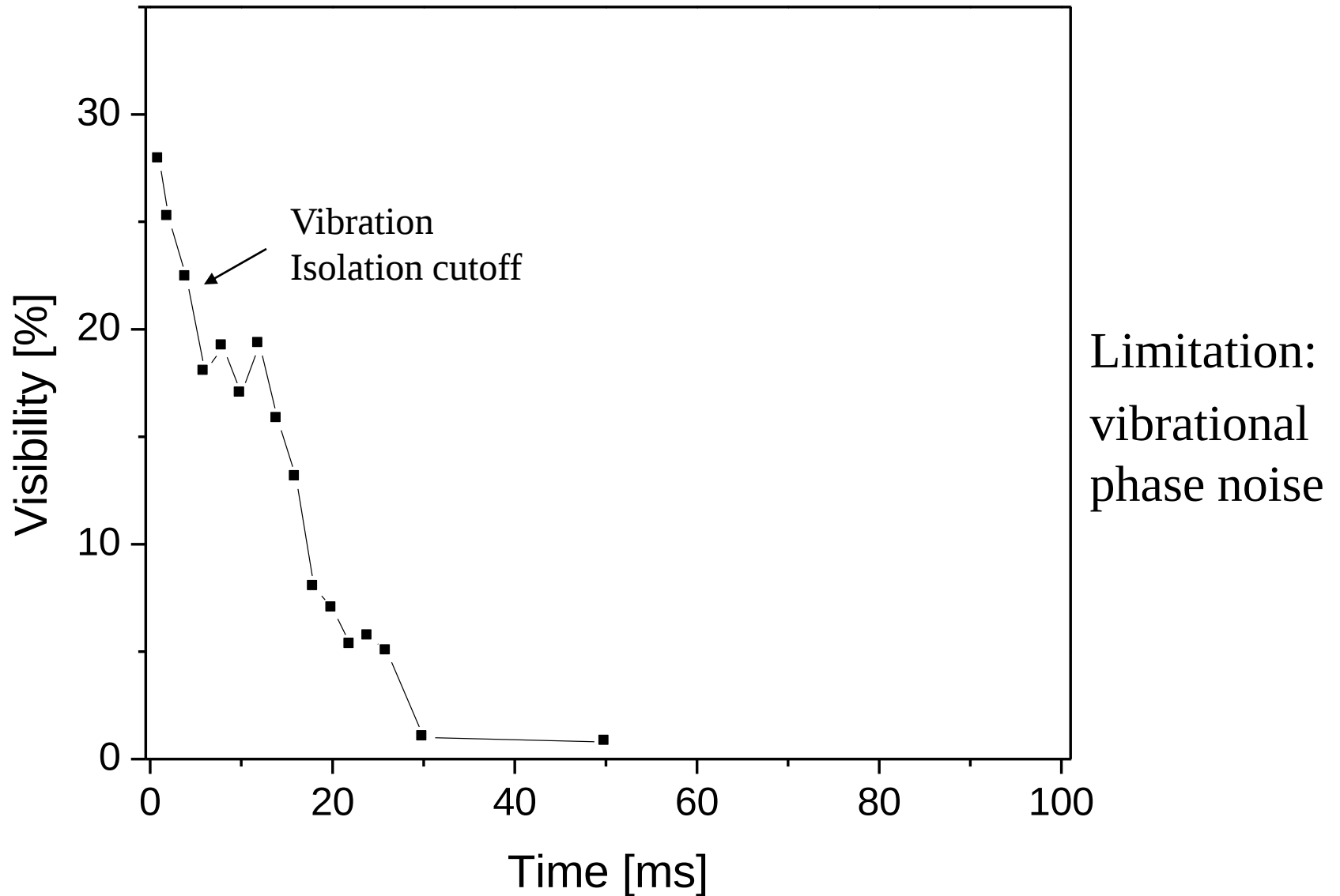


18 photon vs. 2 photon RB Interferometer



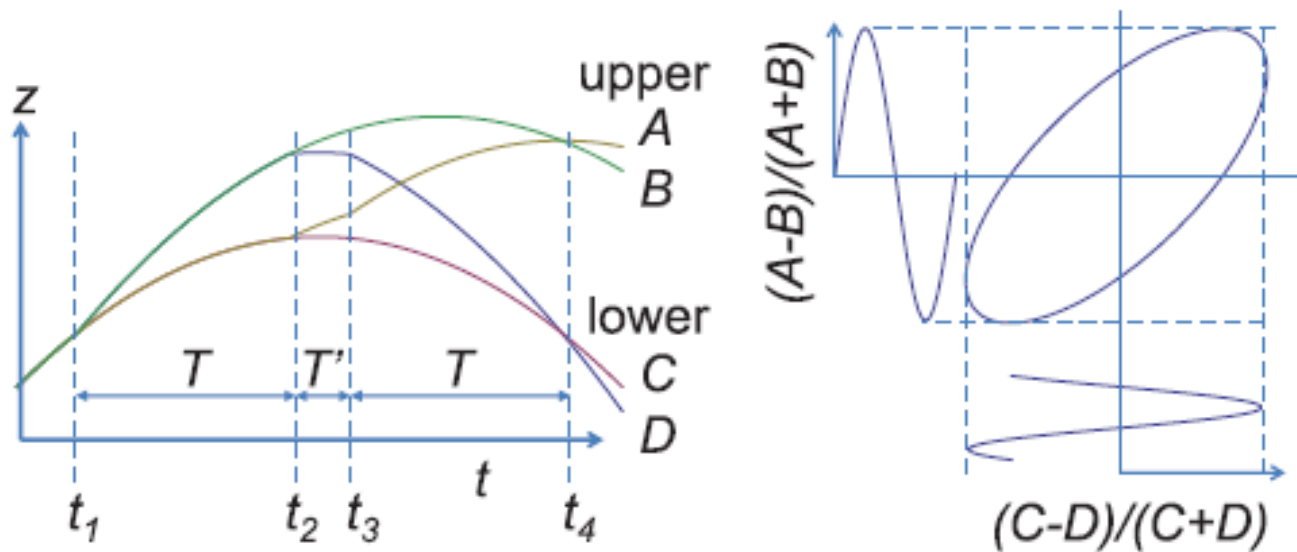


Problem: Contrast decay





Solution: Simultaneous conjugate Interferometers



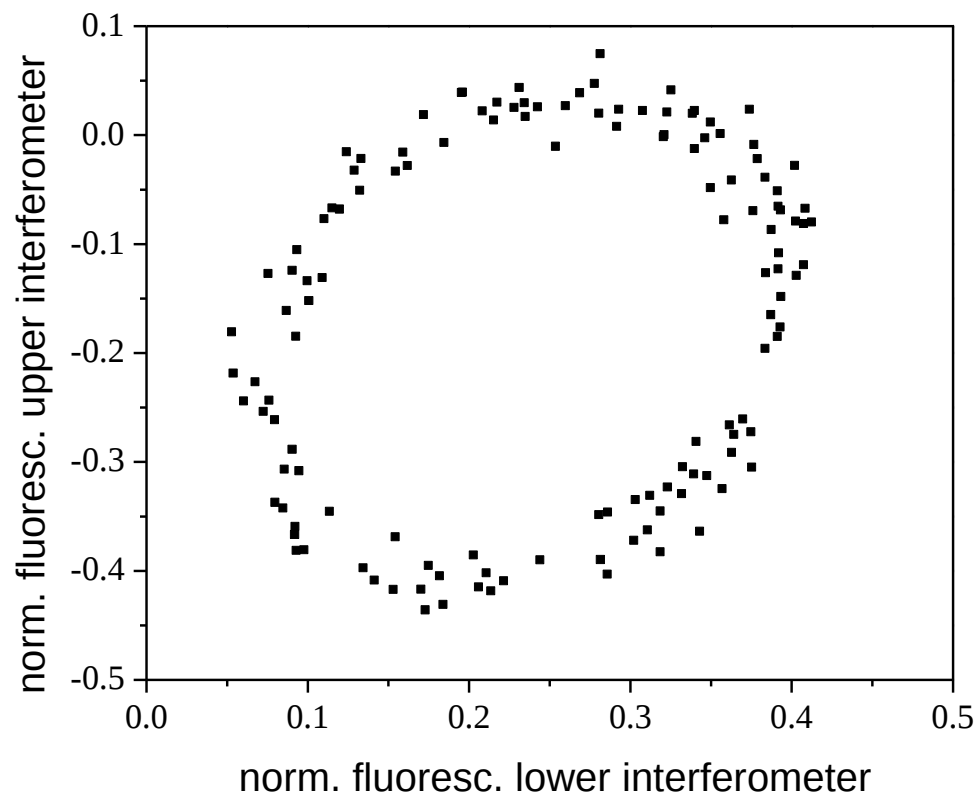
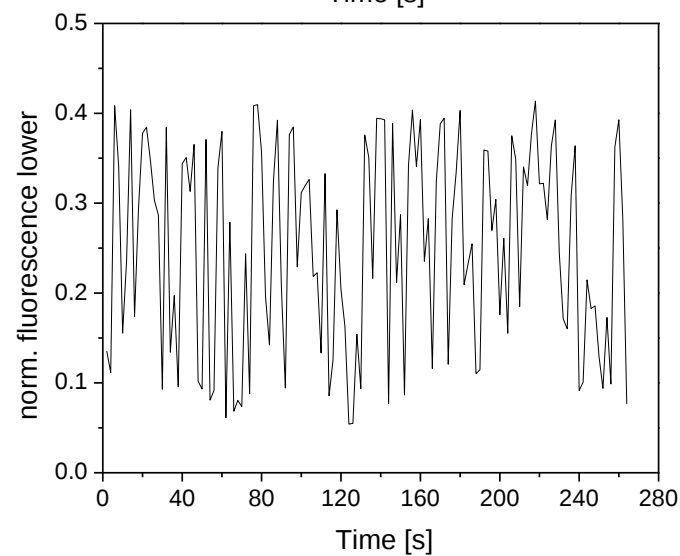
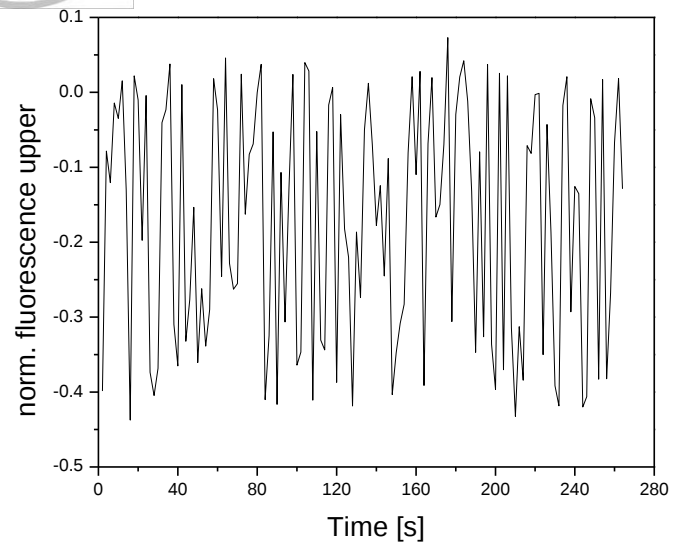
$$\Phi_1 - \Phi_2 = 16\omega_{rec}T$$

Also cancels gravity



Results

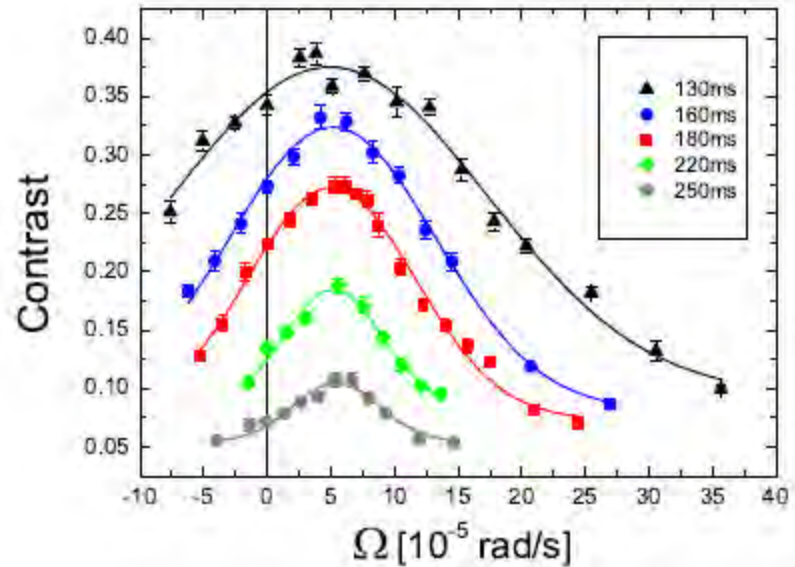
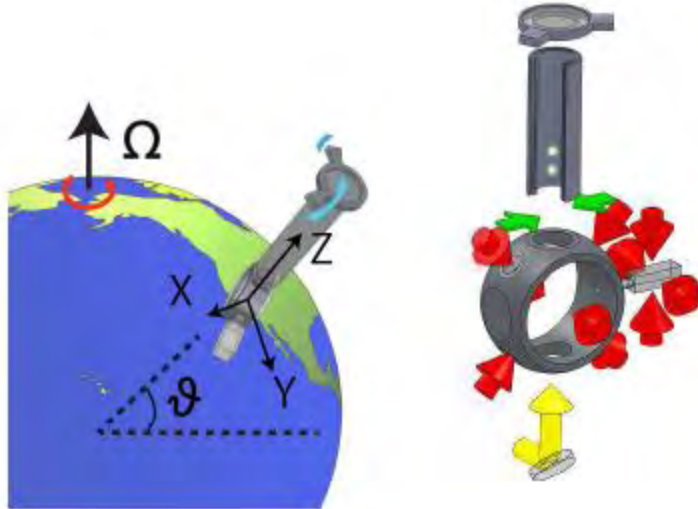
6th order Bragg diffraction, T
=1 ms





Coriolis force

$$\vec{\delta} = 4nv_r\Omega_{\oplus}T(T + T') \cos \vartheta(1, 0, 0).$$



- Rotation causes interferometer not to close
- Can be cancelled by tip tilt-mirror
- Improved contrast (350%), T, and sensitivity
- World's most sensitive atom interferometer (10 ħk, 250 ms)

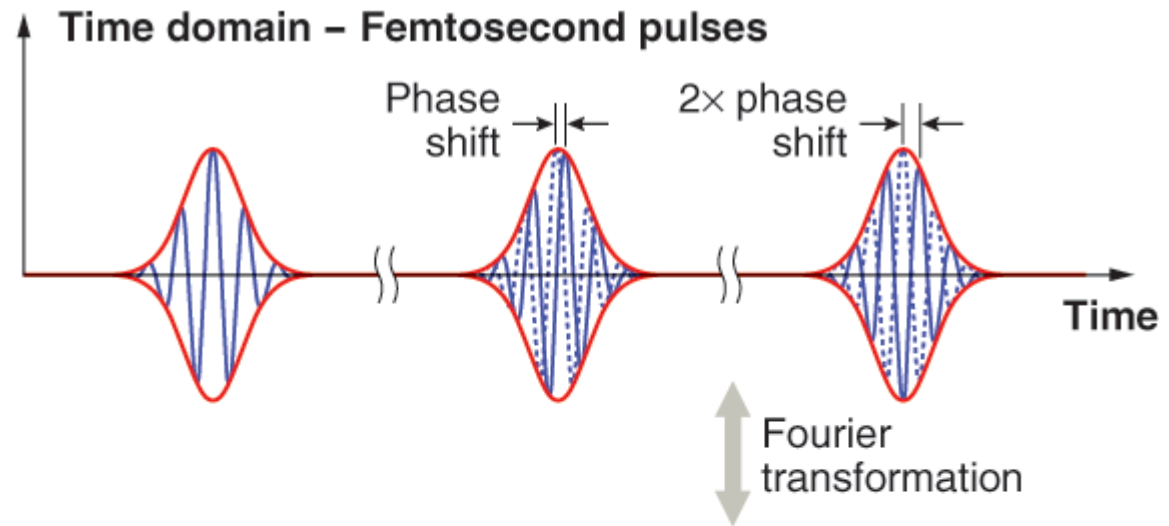
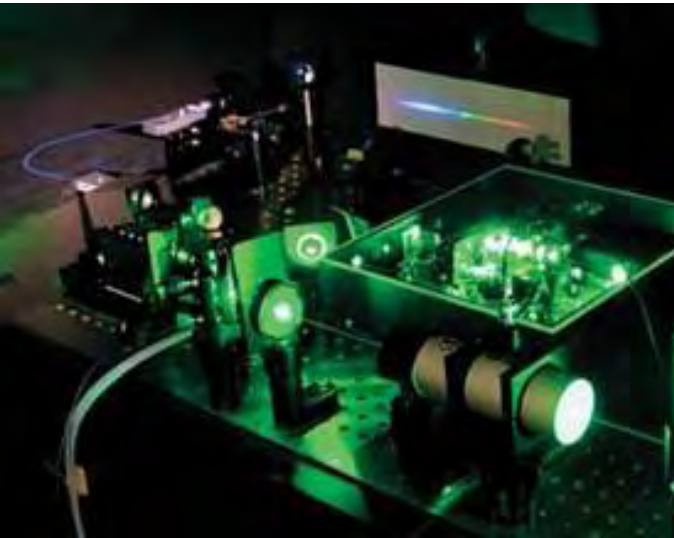


What does this have to do with a
Compton clock?

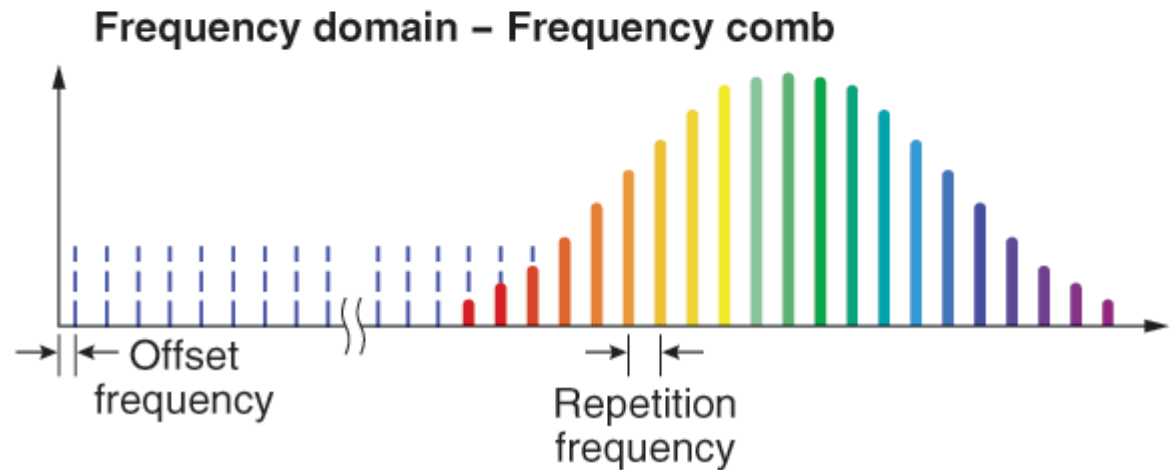
$$\omega = \frac{mc^2}{\hbar} \frac{1}{n^2}$$



Frequency Comb



$$f_n = nf_r + f_{\text{offset}}$$

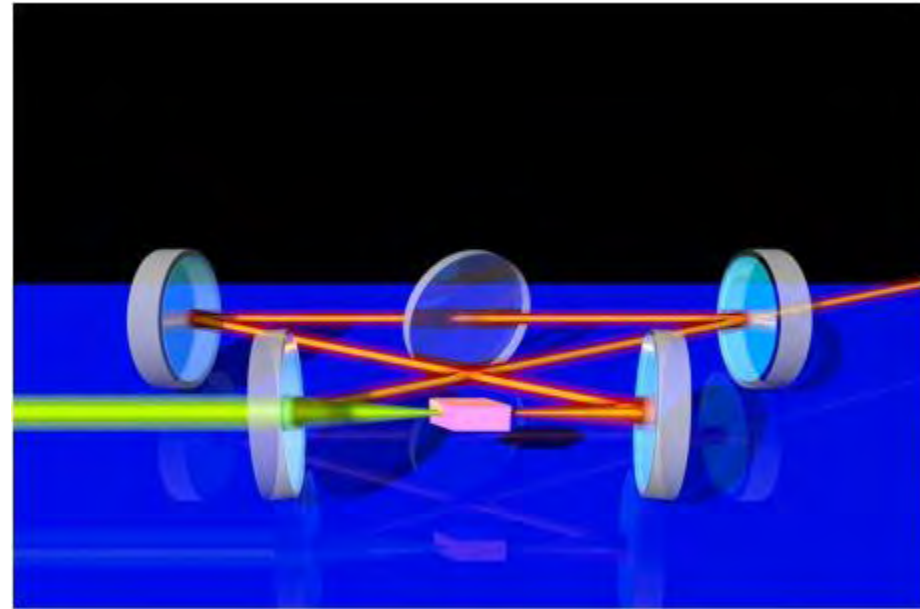


- [1] Theodor Hänsch, Nobel Lecture, http://nobelprize.org/nobel_prizes/physics/laureates/2005/hansch-lecture.html
- [2] J. Hall, Nobel Lecture, http://nobelprize.org/nobel_prizes/physics/laureates/2005/hall-lecture.html

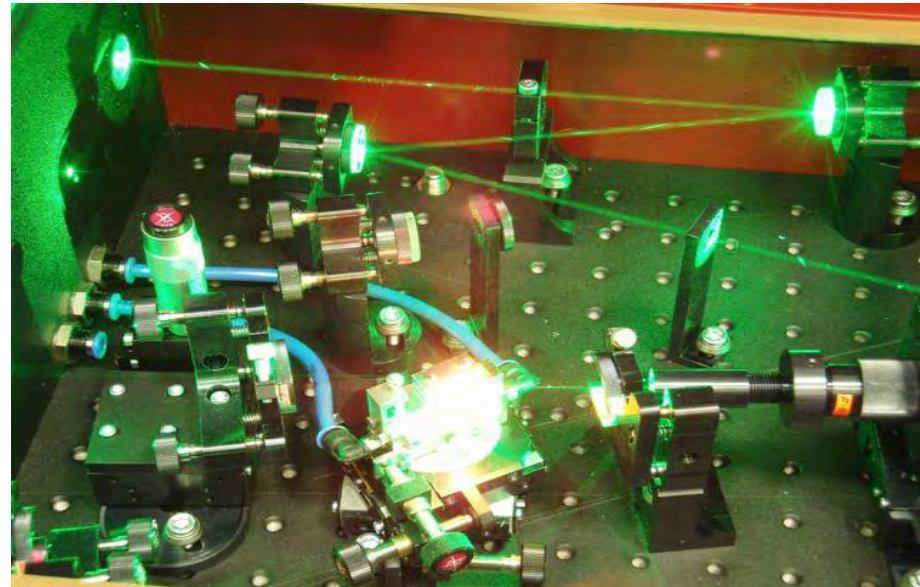


Ti:Sapphire systems

- 500-1000 nm,
- shorter wavelengths by frequency doubling
- 100s of MHz repetition rate

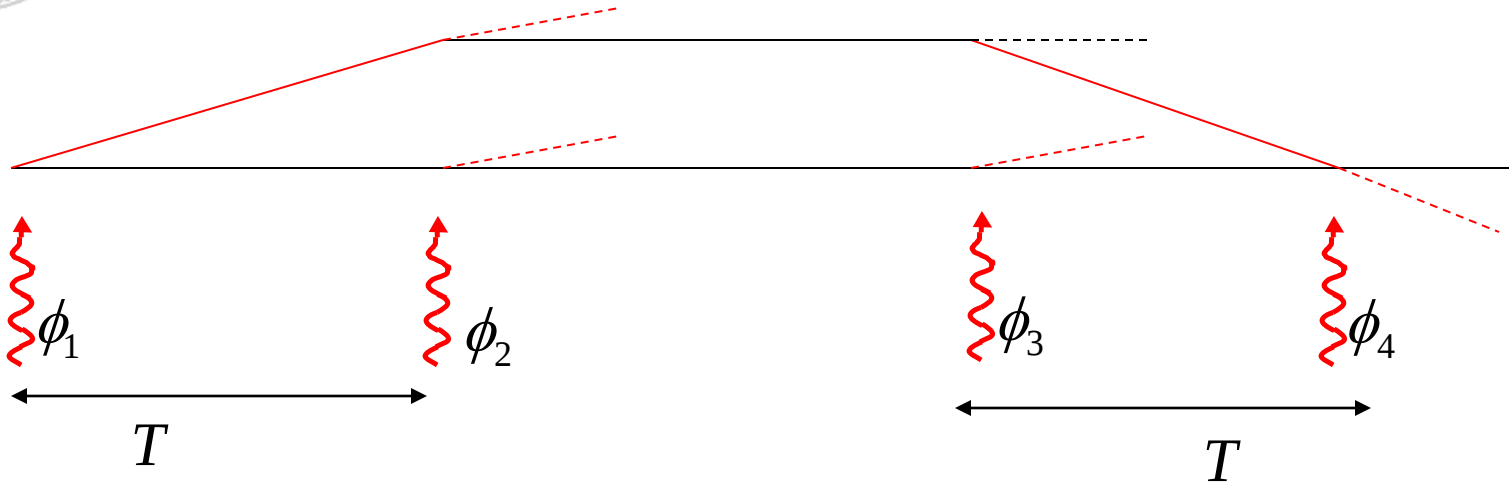


NIST





Ramsey-Borde interferometer



$$\Phi = 8n^2 \frac{\hbar k_L^2}{2m} T = 4n^2 \frac{\omega_L^2}{\omega_C} T$$

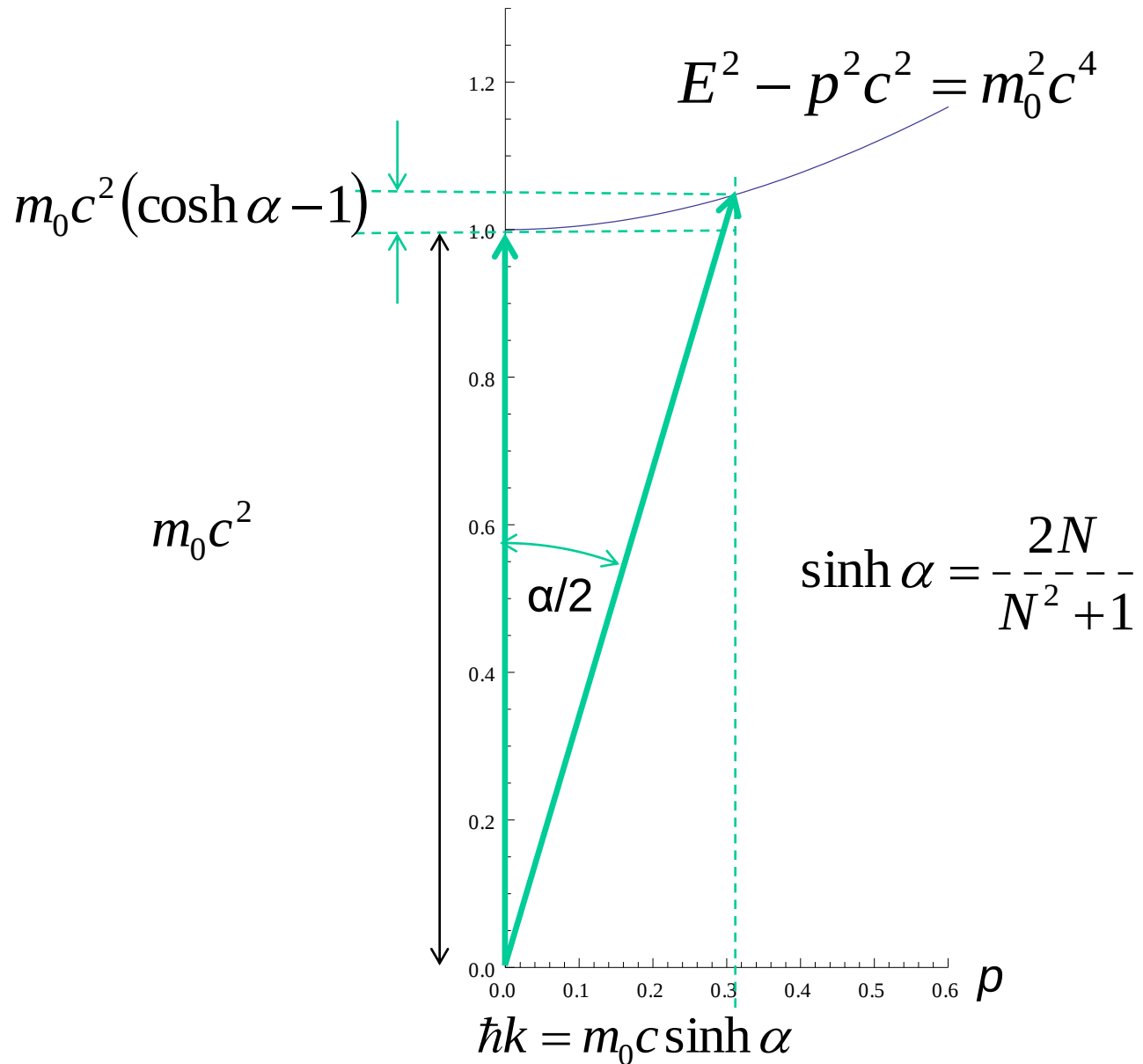
Use frequency comb to stabilize $k = \omega_L / c = N\omega_r$

Then find

$$\Phi_1 = 4n^2 N \omega_C T$$

$$\omega_C : \omega_L : \omega_r = N^2 : N : 1$$

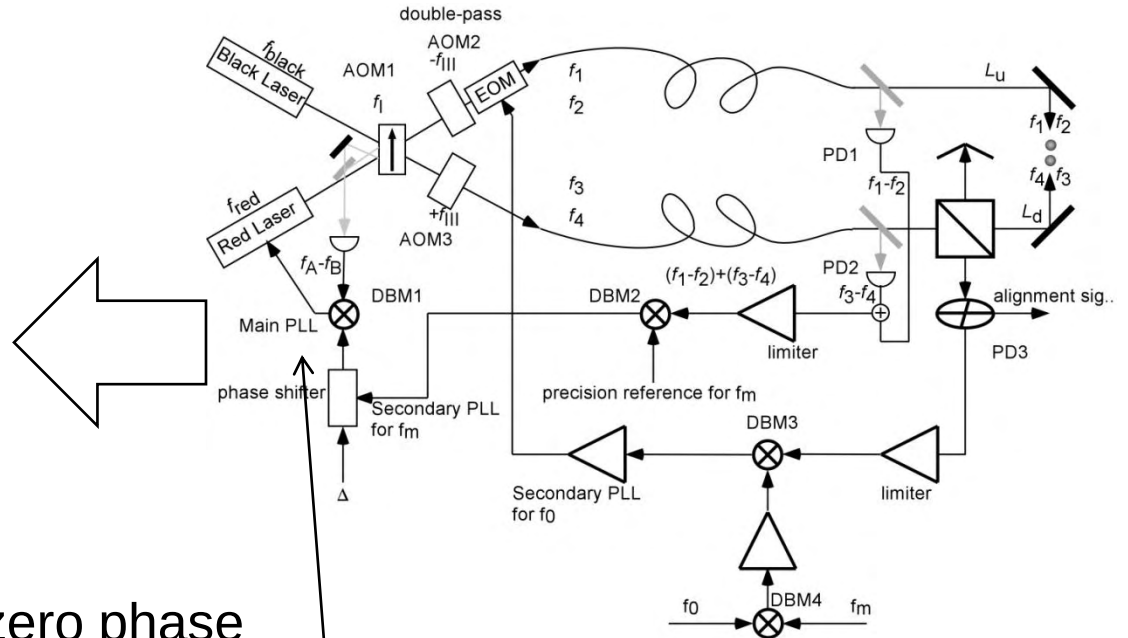
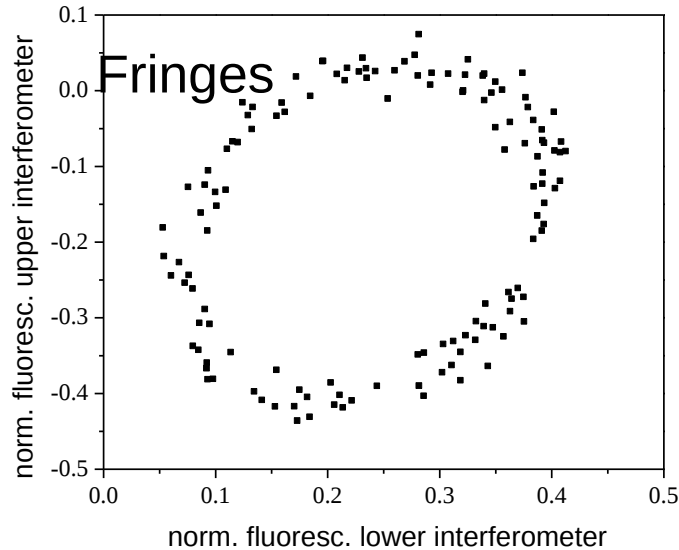
Relativistic derivation





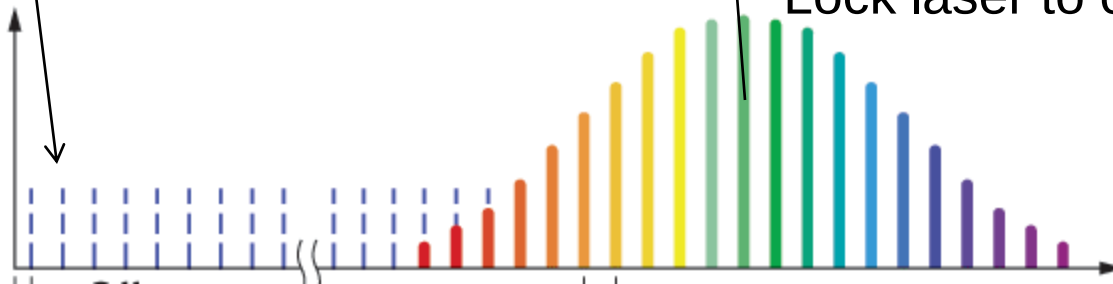
Compton clock

Atom interferometer



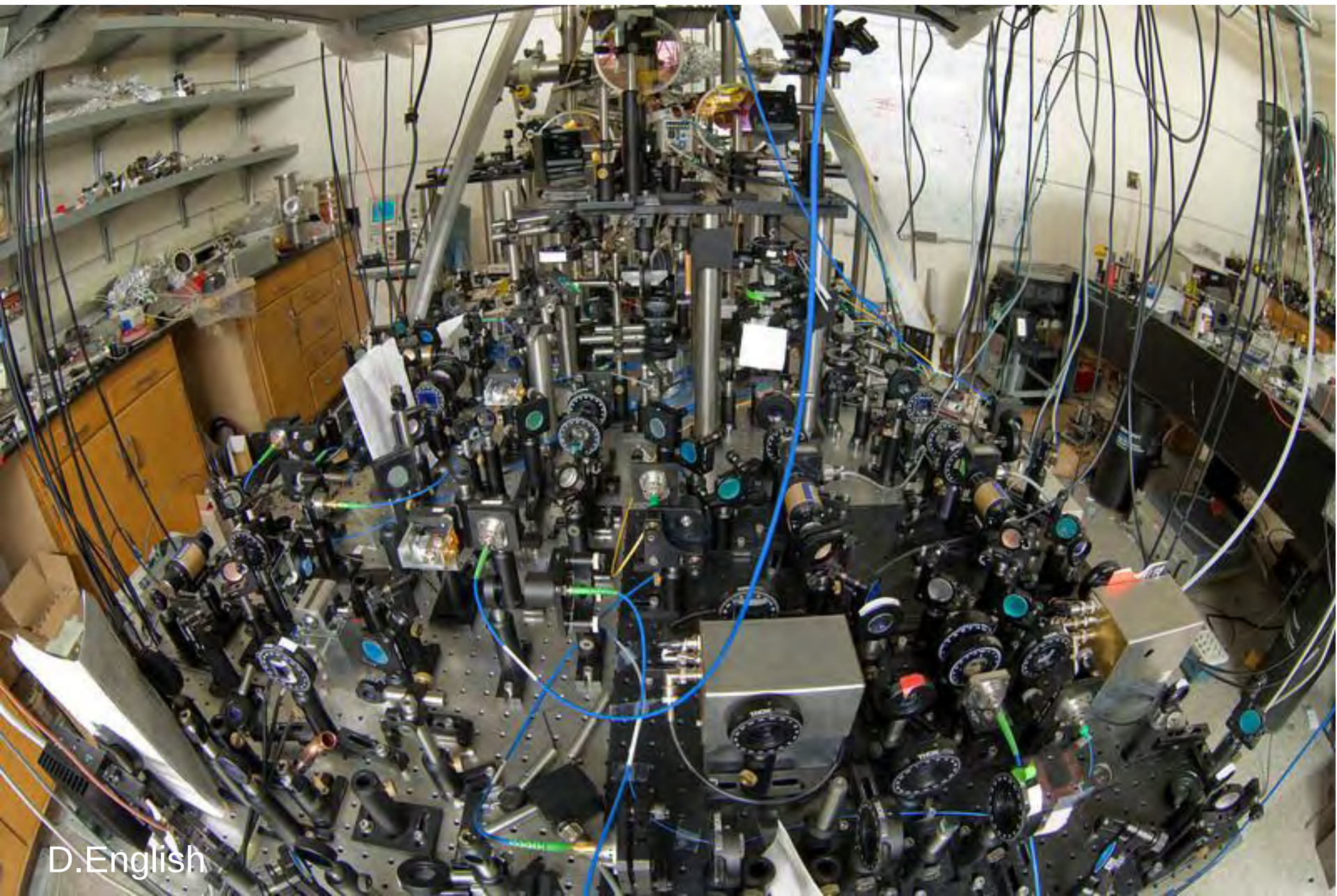
Lock comb to zero phase

Lock laser to comb line

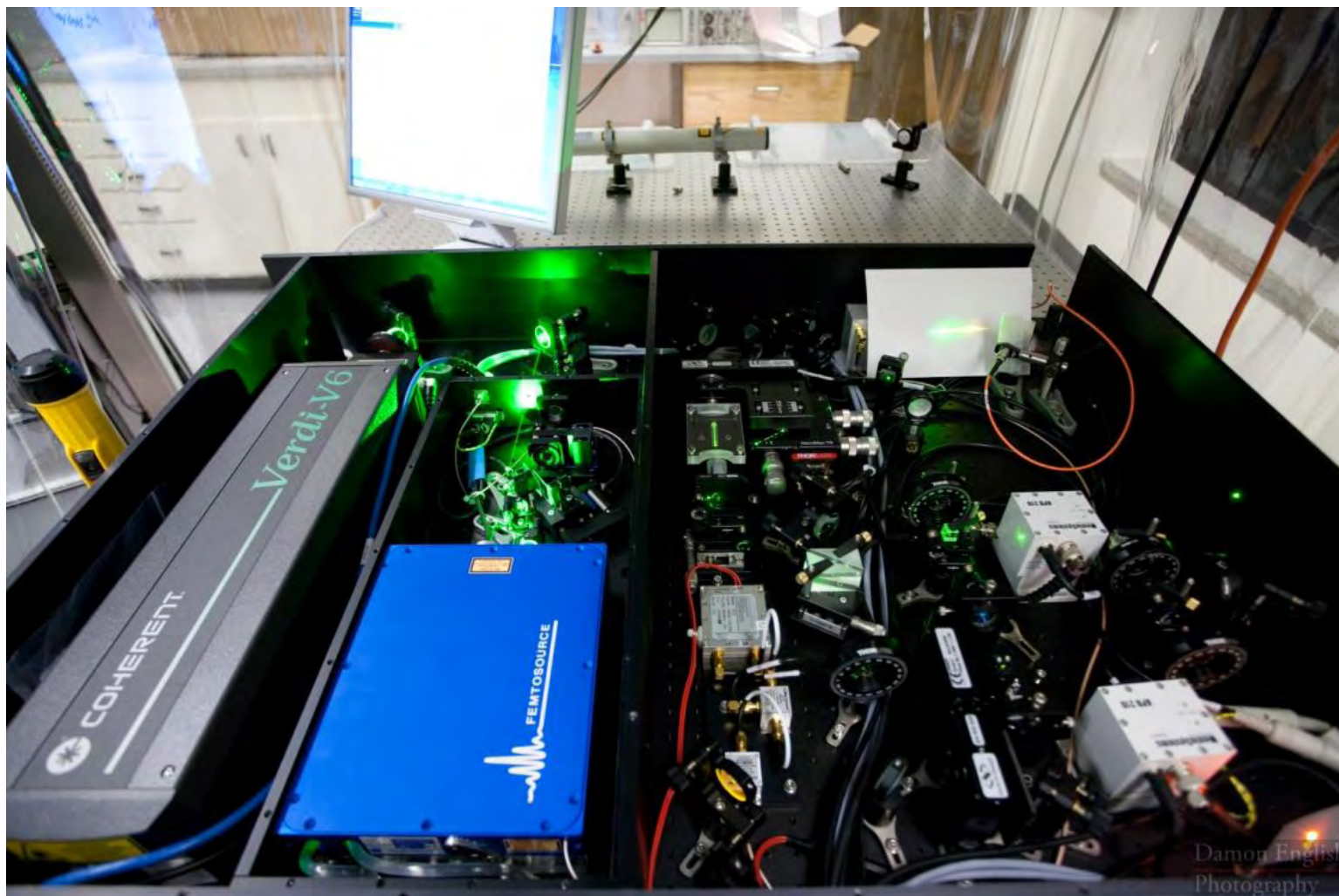


$$\omega_C : \omega_L : \omega_r = N^2 : N : 1$$

3×10^{25} Hz 3×10^{14} Hz 2 kHz

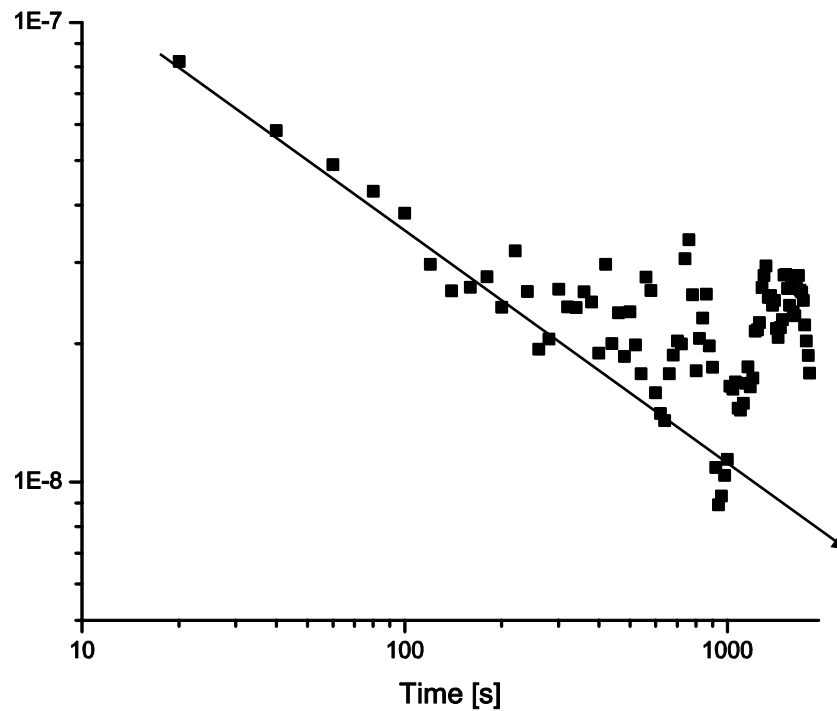
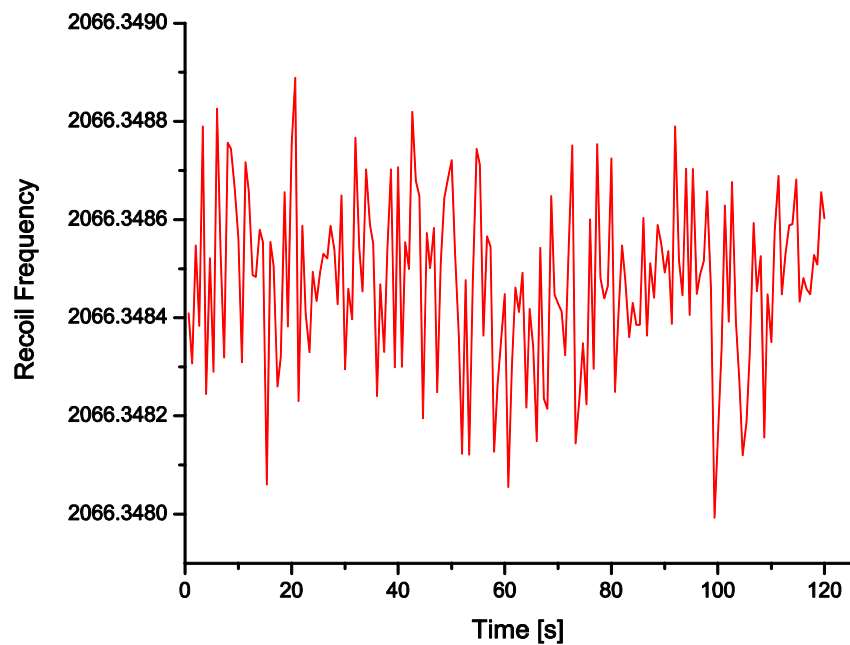


D.English





Preliminary data



Conclusion

- A rock redshifts, time dilates, will soon time the length of the day
- Gravitational AB effect: the phase shift is from the potential, not the force
- Compton clock: divide 3×10^{25} Hz using relativity with self-referenced atom velocity
- 10^{-9} redshift test
- Large momentum transfer, conjugate interferometers

Time is an illusion, lunchtime doubly so

Douglas Adams

$$\Psi = e^{i \frac{mc^2}{\hbar} t} \psi$$

A rock is a clock

Schrodinger equation from GR

$$S = \int mc^2 \sqrt{1 + h_{\mu\nu} u^\mu u^\nu} d\lambda$$

$$\approx \int mc^2 \left(1 + \frac{1}{2} h_{00} + h_{0j} \frac{v^j}{c} - \frac{1}{2} (\delta_{jk} - h_{jk}) \frac{v^j}{c} \frac{v^k}{c} \right) dt$$

$$\Psi(t + \varepsilon, x_A) = N \int d^3 \xi \Psi(t, x_A - \xi) \exp \left(i \frac{mc^2}{\hbar} \left(1 + \frac{1}{2} h_{00} \right) \varepsilon \right) \exp \left(-\frac{1}{2} A_{jk} \xi^j \xi^k + B_j \xi^j \right)$$

$$\int e^{-\frac{1}{2} A_{jk} \xi^j \xi^k + B_j \xi^j} d^3 \xi = \frac{(2\pi)^{3/2}}{\sqrt{\det A}} e^{-\frac{1}{2} B_j (A^{-1})_{jk} B_k}$$

$$i\hbar \frac{d}{dt} \Psi = -mc^2 \frac{1}{2} h_{00} \Psi - \frac{\hbar^2}{2m} (\vec{\nabla} - m\vec{h}) \Psi$$